



power input to that calculated on the $= 3.36 \mu$. Curve a— $m = 3$, $v_e = 0$, A as by Dash (16), A as given by Vinen; curve $m = 4$, $v_e = 0$, $A = 50$ cm-sec/gm.

in the region 1.7° – 2.0° K for large \bar{q} . We have not been able to solve the \bar{q} , but instead we have used a vari- R. B. Lazarus.

$$\frac{\Lambda}{1 + \lambda\delta} d\tau \quad (44)$$

lating α (determined from Vinen's determined from the present experi- fixed and varying λ we obtain

$$\int_{\tau_0}^{\tau} \frac{\Lambda\delta}{(1 + \lambda\delta)^2} d\tau \quad (45)$$

$$\frac{2\lambda\Lambda\delta}{(1 + \lambda\delta)^2} d\tau \left(\frac{\partial \bar{q}}{\partial T} \right)_\lambda \quad (46)$$

$$\frac{(1 + \lambda\delta)L}{\Lambda d^2} \quad (47)$$

where all the quantities are to be evaluated at T_1 . From the experimental data we compute $(\partial T/\partial \bar{q})_\lambda$ and solve (47) for $(\partial T/\partial \lambda)_{\bar{q}}$ with λ set to unity. Since $\Delta T = (\partial T/\partial \lambda)_{\bar{q}} \Delta \lambda$, where ΔT is the temperature difference at (T_1, \bar{q}) between the measured curve and the curve calculated using α (Vinen), we can determine $\Delta \lambda$ and hence $\alpha' = (1 + \Delta \lambda)\alpha$ and new values of A . Table II lists some results of these computations and presents a comparison with values of A as obtained by several other workers. For Slit III' and $T_1 = 1.800^\circ$, 1.900° , and 2.000° K, A has been given for two values of T_0 ; in each case intermediate values of A would be obtained from heating curves beginning at a temperature between these two limits of T_0 . It can be seen from Table II that the experiments are rather well represented by Vinen's values of A . Whereas these considerations may not be useful in any attempt to improve upon Vinen's $A(T)$, it is evident from them that any set of $A(T)$ that is substantially different from those given by Vinen—e.g., as indicated by Brewer and Edwards (17) or by Kramers *et al.* (18)—would not be compatible with the experiments of I and II.

From the arguments presented above and other comparisons with the data of I and II we have concluded that of the various models we have examined, calculations made using Vinen's $A(T)$, $m = 3$, and $v_e = 0$ provide the best overall representation for the experimental data for the 3.36μ and 2.12μ slits. The general character of the agreement may be observed from an examination of Fig. 3 and 4, where families of the heat flow curves for the 3.36μ slit are presented as observed and as computed, respectively. A more quantitative comparison for heat flow is presented in Fig. 5, where $[(\dot{Q}_{\text{obs}} - \dot{Q}_{\text{calc}})/\dot{Q}_{\text{obs}}] \times 100$

TABLE II
COMPARISON OF VALUES OF $A(T)$ OBTAINED FROM HEATING CURVES FOR SLIT III' WITH VALUES OBTAINED BY OTHER WORKERS

$T_1(^{\circ}\text{K})$	$T_0(^{\circ}\text{K})$ (This work only)	$A(T_1)(\text{cm sec gm}^{-1})$			
		Slit III' $d = 3.36 \times 10^{-4}$ cm	Vinen (4) $d = 0.4$ cm, 0.24 cm	Kramers <i>et al.</i> (18) $d = 0.26$ cm	Brewer and Edwards (17) $d = 0.011$ cm, 0.37 cm
1.700	1.083	60 (–7) ^a	75	37 ^b	110
1.800	1.083	98 (6)	91	42	140
	1.586	97 (4)			
1.900	1.083	128 (15)	110	52	185
	1.698	117 (5)			
2.000	1.083	150 (20)	135		260
	1.794	111 (–14)			

^a Numbers in parentheses indicate $(T_{\text{obs}} - T_{\text{calc}})$ in millidegrees at (T_1, \bar{q}) .

^b Note added in proof: In the Proceedings of the Eighth International Conference on Low Temperature Physics (London, England, Sept. 16–22, 1962; to be published) Wiarda and Kramers have reported that new measurements of $A(T)$ are in complete agreement with the results of Vinen.